THE GENERATING FUNCTION TECHNIQUE AMELIORATES EFFECTIVE AND STRING FIELD THEORIES AND FORESHADOWS THEIR LINKAGE TO QUANTUM INFORMATION THEORY

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Abstract

Many scientists are trying to develop a theory of everything or a supposition to explain all aspects of the physical universe. This paper explores a set of theories called effective and string field theory or *EFT* and *SFT*, respectively. These suppositions can be utilized in both old and possibly new physics. Typically, *EFT* and *SFT* have a mathematical method for solving problems called the perturbation theory (*PT*); the generating function technique or GFT can substitute this means of problem-solving. The latter method is used to solve several examples of physical problems, such as determining the cause of muon g-2 experimental deviations, the means for the calculation of glueballs via meson decay, the ascertainment of tetraquark mass from their decay products, and the analysis of binary black hole mergers. Ultimately, *GFT*, instead of traditional *PT* methods, is a potent tool for improving our understanding of concepts in contemporary physics, such as in *EFT* and *SFT*. Also, *GFT* shows the existence of a triality between *EFT*, *SFT*, and Quantum Information theory (*QIT*).

Keywords: generating function technique, effective, quantum, and string field theory.

1.) Introduction

An appropriate theory of everything or *ToE* should adequately combine general relativity and quantum mechanics [2,57,61,62]. Physics significantly grew after Sir Isaac Newton discovered gravity [1,38]. Then, Albert Einstein radically updated humanity's current understanding of gravity with his General Theory of Relativity [13,19,29,37,38,46,91]. General relativity heralded another revolution in physics [1]. Quantum mechanics, the study of quanta or particles, underwent a revolution a little after the world recognized general relativity as being virtually true [18,25,59]. Even though many physicists developed putative mathematical frameworks, there was no provable strong theory combining general relativity and quantum mechanics to date.

Two basic systems of ideas that attempted to explain aspects of the universe involved fields. *EFT* and *SFT* implemented physical body (i.e., particle) and string fields to describe the behavior of facets of classical and quantum physics [21,26,52]. Hypothetically, both theories served as a basis for a new *ToE*.

Calculus was the field of mathematics heavily entrenched in physics; thus, methods for solving such problems incorporating differential equations were essential. Since the inception of calculus, individuals have established various techniques (i.e., perturbation theory or *PT*) for deriving the solutions to differential equations [20]. The latest method for solving physics problems, *GFT*, which this author discovered, involved using several truncated Laurent series of formal power series or generating functions to solve problems known to old physics (i.e., Boussinesq equation, Navier-Stokes problems, etc.) [3,63]. This novel method of accruing solutions to differential equations has a broad reach and is considered capable of solving a wide range of problems in mathematical physics [88].

This paper discusses *EFT* and *SFT* via *PT* as practical *ToEs*. This study is divided into several sections: section 2 provides concepts in perturbation, effective, and string field theories; the next section shows how *GFT* serves as a template for solving problems linked to *EFT* and *SFT*; Section 4 shows how *EFT* and *SFT* via *GFT* can be used in deriving solutions to four issues alluding to old and new physics; finally, the last section gives a quick review of *EFT* and *SFT* that shows *GFT* is a highly effective means for solving problems related to the two suppositions. In addition, the GFT claims at least two significant links are apparent between the two field theories and *QIT* in the last section.

2.) Basics of EFT and SFT after the consideration of PT, then GFT

2.1.) *PT* synopsis

In *PT*, an individual finds an approximate solution to a more straightforward defined problem [64]. The solution becomes more accurate as the approximate solution gains terms with decreasing parameters [64,89]. Ultimately, the approximate solution asymptotically approaches the exact solution as the number of terms added to the perturbation series approaches infinity [64]. In other words, the perturbation series becomes a formal power series over time [64].

2.2.) *EFT* synopsis

EFT encompasses an extensive array of fields in physics [21]. Therefore, *EFT* covers quantum, classical, and cosmological fields [21]. In this study, we will focus on quantum field theory, or QFT, and the cosmological aspects of *EFT*.

QFT combines quantum mechanics, classical field theory, and special relativity [4,24,28,47,60]. It is commonly applied to particle physics and, thus, essential in forming models within subatomic and condensed matter physics [4,24,28,47,60,92]. Since its advent in the 1920s and rebirth in the 1970s, *QFT* has had a prominent role in describing contemporary physics.

QFT was divided into at least three branches: quantum electrodynamics (*QED*), quantum flavor-dynamics (*QFD*), and quantum chromodynamics (*QCD*). *QED* was primarily developed by Dirac in 1927 and was built upon canonical quantization [9,32,40,44]. Also, it dealt with the interaction of fermionic and electromagnetic fields. *QFD* studied electroweak nuclear force, such as bosons Z^0 and $W^{\pm}$ activities, while *QCD* involved nuclear solid interactions, generally mediated gluon fields [12,23,47]. It is expected to find situations where certain branches, like *QED* and *QCD*, cross over or encroach on each other.

PT can be used to solve many problems in *QFT* [47,60,92]. The interaction between particle fields is treated as small perturbations in a free field. Finally, the integration of *PT* with *QFT* is called perturbative quantum field theory or *pQFT*.

EFT can also be applied to cosmology [14,15,29]. Cosmology is the study of the universe and its evolution. This realm of study includes the observations of large-scale structures (*LSS*) and the laws that dictate their behavior [95]. *PT* plays an active role in solving problems associated with cosmology: The combination of *PT* and this area of study is called cosmological perturbation theory [14,15,29]. The practical field theory of large-scale structures or *EFToLSS* is used to enhance the derivation of solutions in this area of science via novel *PT* methodologies [96].

2.3.) *SFT* synopsis

This supposition involves reformulating relativistic strings to *QFT* [12,24,28,60,92]. Unlike *QFT*, which treats an individual field of particles in excited states or quanta, string theory converts point-like particles into one-dimensional entities, referred to as strings. Thus, string field theory uses these one-dimensional strings to define excited particle fields.

Second quantization dictates the type of SFT to be considered, such as open and closed string fields [97]. In the standard model of particle physics, some open string fields are represented by gauge (gluons, photons, W and Z bosons) and quark or lepton fermions [90]. Fields that describe the scattering of open or closed strings are called open or closed SFT, respectively. However, if a field contains a combination of both open and closed strings, it is referred to as an open-closed SFT.

SFT possesses advantages over regular string theory. For instance, it permits the calculation of "off-shell" amplitudes and thus provides information about string scattering [52]. In addition to giving an individual the means to calculate the masses of particle systems that obey classical equations of motion, it can be used to determine particle systems that do the contrary. This process is called "off the mass shell" or off-shell of the mass hyperboloid via perturbation methods [52]. In other words, SFT can be used to define attributes of particles that do not follow the equations of motion in a classical sense (i.e., virtual particles, dark sector quanta, etc.) due to its innate ability to depict particles as string fields.

2.4.) Mathematical descriptions of effective and string fields via *GFT*

In *EFT*, an elementary field γ is a formal power series of a function f. A physical body, such as a particle, can be designated as an exponential function f, assuming the auxiliary/characteristic equation is of the first order, takes the following form:

$$f(\xi) = c_1 e^{-\xi},$$

where c_1 is an arbitrary constant, and ξ is the ansatz transform variable. (*Note: For GFT purposes, the auxiliary function is generally designated by the Greek letter ϕ. However, this study will designate the auxiliary function with the alphanumeric letter f due to particle physicists' conventional use of ϕ for bosons.*) The ansatz transformed variable ξ for a (3+1) system was defined as

$$\xi = \alpha t + \beta_1 x + \beta_2 y + \beta_3 z.$$

On the other hand, the expression of a "brane," which can be designated by a sinusoidal wave function f, Assuming the auxiliary/characteristic equation is of the second order and is expressed as

$$f(\xi) = c_1 \cos(\xi) + c_2 \sin(\xi)$$

or

$$f(\xi) = c_1 \cosh(\eta) + ic_2 \sinh(\eta)$$

where η is the **imaginary** ansatz transform variable, where $\eta = i\xi$, and c_2 is another arbitrary constant.

The physical body and wave function f leads to *EFT* and *SFT*, respectively.

Thus, the elementary physical body or string field γ can be defined as

$$\gamma(\xi) = \sum_{k=0}^{\infty} p_k f(\xi)^k,$$

where p_k is the *k*-th parameter/coefficient of the formal power series γ. On the contrary, its conjugate elementary qubit gate γ^* is the following expression:

$$\gamma^*(\xi) = -\sum_{k=0}^{\infty} p_k f(-\xi)^k.$$

The formal power series g and $g^*$are also elementary effective or string fields. It is important to state an elementary string field is an object with an infinite number of branes; one of the branes acts as the "string" while the other branes serve as the "bulk" space or compactified dimensions of the universe [82,87]. Finally, a truncated Laurent series of the elementary effective or string field g or g^* raised by some power j forms a transformed compound effective or string field U if p_k is combinatorial or trigonometric:

$$U(\xi) = \sum_{j=-n}^{n} q_j \gamma(\xi)^j$$

or

$$U'(\xi) = -\sum_{j=-n}^{n} q_j \gamma'(-\xi)^j,$$

where n is the absolute integer value of the truncated power and q_j is the j-th parameter/coefficient and power. Regarding string field theory, n is equal to the supersymmetry level N. Ultimately, the difference between PT and GFT is that one method builds upon an approximate solution while the other narrows the general solution to derive the exact solution. The former method would require many steps to achieve its objective due to adding higher-order terms. At the same time, the latter only needs a few steps, like solving the parameter/coefficient and arbitrary constants.

3.) GFT as a new mathematical basis for garnering solutions in EFT or SFT

3.1.) *GFT* and general solutions to particles

GFT, with some modification, was a method implemented to find the solution of [non]linear PDEs. Its transformed general solution *U* comprised a Laurent series set of combinatorial or trigonometric-based generating functions [3]. One should consider the transformed general solution *U* as a transformed effective or string field. After one assessed the maximal and minimal power n, through which the Laurent series of various types of elementary practical or string field γ or γ', their conjugate is eventually truncated, (s)he plugs in the predefined function f into the transformed general solutions for effective or string fields ϕ_l, ϕ_l', ϕ_l^{μ}, $\phi_l'^{\mu}$, $\Phi_l^{\mu\nu}$, $\Phi_l'^{\mu\nu}$, Ψ_m, and Ψ_m':

$$\Phi_l(\xi) = \sum_{i=1}^{2}\sum_{j=-\mathrm{n}}^{\mathrm{n}}(al_{ij}(\sum_{k=0}^{\infty}2f(\xi)^k S_k(0)^i)^j + bl_{ij}(\sum_{k=0}^{\infty}2C_k(0)^i f(\xi)^k)^j),$$

$$\Phi_l'(\xi) = -\sum_{i=1}^{2}\sum_{j=-\mathrm{n}}^{\mathrm{n}}(al_{ij}(\sum_{k=0}^{\infty}2f(-\xi)^k S_k(0)^i)^j + bl_{ij}(\sum_{k=0}^{\infty}2C_k(0)^i f(-\xi)^k)^j),$$

$$\Phi_l^{\mu}(\xi) = e^{-\tau}\sum_{i=1}^{2}\sum_{j=-\mathrm{n}}^{\mathrm{n}}(al_{ij}^{\mu}(\sum_{k=0}^{\infty}2f(\xi)^k S_k(0)^i)^j + bl_{ij}^{\mu}(\sum_{k=0}^{\infty}2C_k(0)^i f(\xi)^k)^j),$$

$$\Phi_l'^{\mu}(\xi) = -e^{\tau}\sum_{i=1}^{2}\sum_{j=-n}^{n}(al_{ij}'^{\mu}(\sum_{k=0}^{\infty}2f(-\xi)^k S_k(0)^i)^j + bl_{ij}'^{\mu}(\sum_{k=0}^{\infty}2C_k(0)^i f(-\xi)^k)^j),$$

$$\Phi_l^{\mu\nu}(\xi) = \sum_{i=1}^{2}\sum_{j=-\mathrm{n}}^{\mathrm{n}}(al_{ij}^{\mu\nu}(\sum_{k=0}^{\infty}2f(\xi)^k S_k(0)^i)^j + bl_{ij}^{\mu\nu}(\sum_{k=0}^{\infty}2C_k(0)^i f(\xi)^k)^j),$$

$$\Phi_l'^{\mu\nu}(\xi) = -\sum_{i=1}^{2}\sum_{j=-\mathrm{n}}^{n}(al_{ij}'^{\mu\nu}(\sum_{k=0}^{\infty}2f(-\xi)^k S_k(0)^i)^j + bl_{ij}'^{\mu\nu}(\sum_{k=0}^{\infty}2C_k(0)^i f(-\xi)^k)^j),$$

$$\Psi_m(\xi) = e^{-\tau}\sum_{i=1}^{2}\sum_{j=-\mathrm{n}}^{\mathrm{n}}(cm_{ij}(\sum_{k=0}^{\infty}2f(\xi)^k S_k(0)^i)^j + dm_{ij}(\sum_{k=0}^{\infty}2\mathrm{C}_k(0)^i f(\xi)^k)^j),$$

and

$$\Psi_m'(\xi) = -e^{\tau}\sum_{i=1}^{2}\sum_{j=-\mathrm{n}}^{\mathrm{n}}(cm_{ij}(\sum_{k=0}^{\infty}2f(-\xi)^k S_k(0)^i)^j + dm_{ij}(\sum_{k=0}^{\infty}2\mathrm{C}_k(0)^i f(-\xi)^k)^j),$$

where Φ_l and Φ_l' are the transformed scalar bosonic effective or string fields, $\Phi_l^{\mu}$and $\Phi_l'^{\mu}$ are transformed vector bosonic effective or string fields, $\Phi_l^{\mu\nu}$ and $\Phi_l'^{\mu\nu}$ are the transformed tensor bosonic effective or string fields, Ψ_m and Ψ_m^* are the transformed fermionic effective or string fields, and τ is a specific ansatz transformed time variable or

$$\tau = \alpha_\tau\, t.$$

The parameter/coefficient p_k was defined as combinatorial numbers to the *i*-th power like

$$p_k = 2\mathrm{S}_k(0)^i$$

and

$$p_k = 2\mathrm{C}_k(0)^i.$$

The square root of the *k*-th Fibonacci number at/about zero, or $S_k(0)$, was

$$S_k(0) = \sin\left(\frac{\pi k}{2}\right),$$

while the *k*-th Chebyshev *T* or *U* (not to be confused with the transformed general solution of the effective or string field) number at/about zero, or $C_k(0)$, was

$$C_k(0) = \cos\left(\frac{\pi k}{2}\right).$$

For this article, the parameter/coefficient q_j was either al_{ij} or bl_{ij} for bosonic effective or string field while the parameter/coefficient q_j was either cm_{ij} or dm_{ij} were used for the fermionic effective or string fields, where $l = 1, 2, .., n_l$ and $m = 1, 2, .., n_m$.

Ultimately, mesonic effective or string fields exhibit as a hyperbolic secant function raised by some power. In contrast, gauge bosonic and fermionic effective or string fields can be expressed as a logistic function raised by some power. Figure 1 claims odd integer spin bosonic or half-spin fermionic string fields possess open strings whose vibratory modes satisfy solely Dirichlet boundary conditions; the strings, called D-branes, about such fields, have endpoints that are fixed in spacetime [83,86,87]. Thus, *GFT* implies the odd integer spin bosonic or half-spin fermionic string field's cm_{ij} are equal to *null*. On the contrary, Figure 1 suggests that other string fields, such as fields that comprise even integer spin bosonic strings (excluding those that describe scalar particles), can be closed or open strings [76,85,86,87]. They possess vibratory modes that satisfy solely Neumann boundary conditions; these fields have string endpoints that are free to roam in spacetime [76,85,86,87]. In short, open string fields that behave like gauge bosons, such as photons and gluons, have al_{ij}^{μ} equal to *null* and bl_{ij}^{μ} exist as a 4 X 1 vector, while open string fields that behave like pseudoscalar and vector mesons retain a bl_{ij} equal to *null*. On the contrary, closed string fields that behave like tensor bosons possess $bl_{ij}^{\mu v}$ equal to *null* and $al_{ij}^{\mu v}$ exist as a 4 X 4 matrix. The endpoints of the open string fields are fermions.

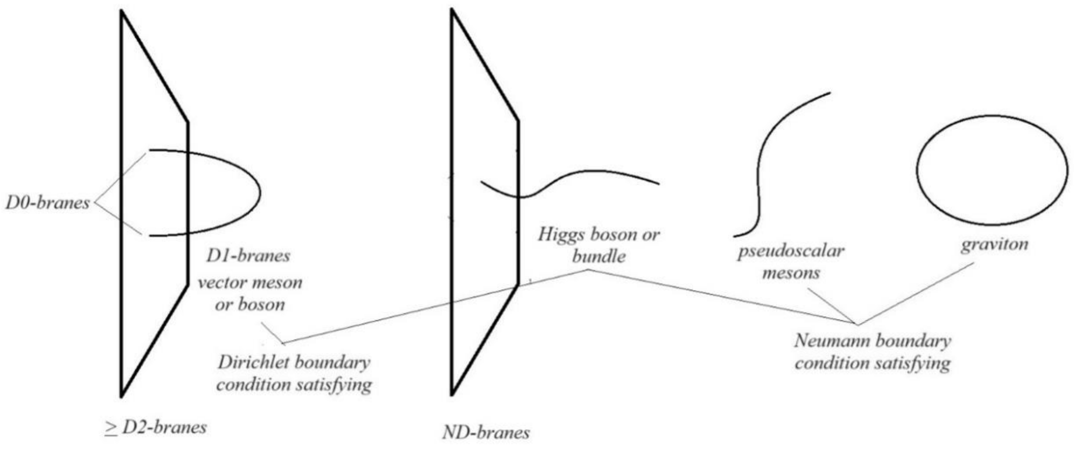


Figure 1: examples of relativistic strings.

For this paper's purposes, al_{ij}^{μ}, bl_{ij}^{μ}, and $al_{ij}^{\mu v}$ can be expressed as

$$al_{ij}^{\mu} = al_{ij}(1,1,1,1),$$
$$bl_{ij}^{\mu} = bl_{ij}(1,1,1,1),$$

and

$$al_{ij}^{\mu v} = al_{ij}\begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}.$$

Finally, most operations involve Hadamard multiplication, mainly if limited to two vectors.

3.2.) On- and off-shell rest mass assessment via renormalization

The expression for renormalization within a volume V was:

$$m_U = \frac{1}{2}\int |U(\xi)U^*(\xi)|dV,$$

where m_U was the mass-energy equivalence for a compound effective or string field U and its conjugate U^*. Assuming the spherical volume, using Manhattan/taxicab-like distance ξ, for compound effective or string field U was equal to the following expression [75]:

$$V = \frac{\pi\xi^3}{6},$$

the formula for renormalization became the following:

$$m_U = |2\int_0^{\infty}\frac{1}{4}\pi\xi^2 U(\xi)U^*(\xi)d\xi|,$$

or

$$m_U = |\int_{-\infty}^{\infty}\frac{1}{4}\pi\xi^2 U(\xi)U^*(\xi)d\xi|.$$

If the field pertains to either a vector or tensor boson, an individual must find the (root) mean (square) of m_U for the former and implement a trace on m_U, then divided by four to the latter.

4.) Examples

This section explores and expands upon two basic Lagrangian equations needed to solve the effective and string fields associated with this paper. One Lagrangian density formalism for either a fermion or gauge boson field involving the Dirac equation, with its complex conjugate field:

$$L[\psi] = h(\partial_\mu\psi\,\partial^\mu\psi' + g_2|\psi|^2)$$

or

$$L[\phi^\mu] = h(\phi^\mu\,\partial^\mu\phi'^\mu + g_2|\phi^\mu|^2).$$

Either equation obeys canonical commutation relations and is converted to the Schrodinger equation if the process is Markovian or the quantum telegraph equation if the process is not Markovian [105]. For this paper's intentions, all functional derivatives of the above equations will not involve a Markov process, yielding the quantum telegraph equation. On the other hand, the inhomogeneous Klein-Gordon equations for a (pseudo)scalar, vector, and tensor mesons/boson are

$$L[\phi] = \tfrac{1}{2}\partial_\mu\phi\,\partial^\mu\phi' - \tfrac{1}{2}c^2M^2\phi\phi' + L_{interactions},$$
$$L[\phi^\mu] = \tfrac{1}{2}\partial_\mu\phi^\mu\,\partial^\mu\phi'^\mu - \tfrac{1}{2}c^2M^2\phi^\mu\phi'^\mu + L_{interactions},$$

and

$$L[\phi^{\mu v}] = \tfrac{1}{2}\partial_\mu\phi^{\mu v}\,\partial_v\phi'^{\mu v} + \tfrac{1}{2}\partial_v\phi^{\mu v}\,\partial_\mu\phi'^{\mu v} + \tfrac{1}{2}\partial_\mu\phi^{\mu v}\,\partial^\mu\phi'^{\mu v} - \tfrac{1}{2}c^2M^2\phi^{\mu v}\phi'^{\mu v} + L_{interactions},$$

respectively. (*Note: If a gauge boson is a constituent of another particle, one should use the quantum telegraph Lagrangian formalism. On the contrary, if a vector meson is considered a composite particle, then (s)he should use the inhomogeneous Klein-Gordon Lagrangian equation.)* The values g_2 is an interaction constant generally equal to unity, while the value h equals a positive integer if the particle/string is presented initially or a negative integer if the particle/string is a by-product [16,44]. Also, the number h equals the coefficient associated with a specific term in the inhomogeneous Klein-Gordon Lagrangian. The principles of least action have several bold Lagrangian terms: ϕ represents the set of scalar bosonic effective or string fields, ϕ^{μ} represents the set of vector-based bosonic effective or string fields, $\phi^{\mu\nu}$ represents the set of tensor bosonic effective or string fields. On the other hand, ψ_i and ψ_f signify the initial and final fermionic effective or string fields, respectively. Also, $'$ symbolizes the conjugate fields.

Supplementary material included with this study is Mathematica® spreadsheets of the following examples of *EFT* solved by *GFT*. The supplementary material also contains an instance of linearized gravitational waves solved via GFT. **Finally, the results between *EFT* and *SFT* were equal!**

4.1.) The cause of the deviation in the muon g-2 experiment

The muon g-2 experiment attempts to measure a muon's magnetic dipole moment or g-factor accurately, and Fermilab is currently conducting it [17,55]. The experiment involves injecting the muons from the decay of pseudoscalar pions into a storage ring and then measuring the muon's g-factor. Theoretically, the strength of the magnetic dipole moment is supposed to be exactly two. Any deviation from this, the latter value claims there are likely additional particles in the standard model in the storage ring.

So far, the laboratory has experimented two times. The results of the first two experiments suggested that the anomalous magnetic moment, which is $a_{\mu} = \frac{g-2}{2}$, had deviated by a factor of 0.00116592 [33]. Ultimately, this data implied there was likely at least one additional particle that is a by-product of a pion but not currently listed in the standard model.

Assuming the additional particle is a vector boson ϕ_f^{μ} and the Lagrangian term for interactions is as follows:

$$L_{interactions}[\phi_i, \phi_f^{\mu}, \psi_{f1}, \psi_{f2}] = g_1\phi_i\left(\left|\psi_{f1}\right|^2 + \left|\psi_{f2}\right|^2 + 2\left|\phi_f^{\mu}\right|^2\right)$$

then the principle of least action for pion decay is as defined as:

$$S[\phi_i, \phi_f^{\mu}, \psi_{f1}, \psi_{f2}] = \int dx^4\left(\begin{array}{c} -g_1\phi_i\left(\left|\psi_{f1}\right|^2 + \left|\psi_{f2}\right|^2 + 2\left|\phi_f^{\mu}\right|^2\right) - g_2\left(\left|\psi_{f1}\right|^2 + \left|\psi_{f2}\right|^2 + 2\left|\phi_f^{\mu}\right|^2\right) \\ -2\,\partial_{\mu}\phi_f^{\mu}\,\partial^{\mu}\phi_f'^{\mu} - \partial_{\mu}\psi_{f1}\,\partial^{\mu}\psi_{f1}' - \partial_{\mu}\psi_{f2}\,\partial^{\mu}\psi_{f2}' \\ -\frac{1}{2}c^2M^2\phi_i\phi_i' + \frac{1}{2}\partial_{\mu}\phi_i\,\partial^{\mu}\phi_i' \end{array}\right).$$

Claiming ϕ_i is the pionic effective or string field, ϕ_f^{μ} is the unknown bosonic effective or string field, and ψ_{f1} is one of the valence muonic effective or string fields, while ψ_{f2} is the valence muon neutrino-based effective or string field associated with the decay of the pion, the functional derivative of the principle of least action yields at least four transformed Hamiltonian equations:

$$\alpha i\Psi_{f1\xi} + \alpha_{\tau} i\Psi_{f1\tau} - \frac{\alpha_{\tau}^2}{c^2}\Psi_{f1\tau\tau} - \frac{\alpha\alpha_{\tau}}{c^2}\Psi_{f1\xi\tau} - \left(\frac{\alpha^2}{c^2} - (\beta_1^2 + \beta_2^2 + \beta_3^2)\right)\Psi_{f1\xi\xi} = g_1\Phi_i\Psi_{f1} + g_2\Psi_{f1},$$

$$\alpha i\Psi_{f2\xi} + \alpha_{\tau} i\Psi_{f2\tau} - \alpha_{\tau}^2/c^2\Psi_{f2\tau\tau} - \alpha\alpha_{\tau}/c^2\Psi_{f2\xi\tau} - \left(\alpha^2/c^2 - (\beta_1^2 + \beta_2^2 + \beta_3^2)\right)\Psi_{f2\xi\xi} = g_1\Phi_i\Psi_{f2} + g_2\Psi_{f2},$$

$$\left(\alpha^2/c^2 - (\beta_1^2 + \beta_2^2 + \beta_3^2)\right)\Phi_{i\xi\xi} + c^2M^2\Phi_i = g_1\left(\Psi_{f1}\Psi_{f1}' + \Psi_{f2}\Psi_{f2}' + 2\Phi_f^{\mu}\Phi_f'^{\mu}\right),$$

and

$$\alpha i\Phi_{f\xi}^{\mu} + \alpha_{\tau} i\Phi_{f\tau}^{\mu} - \alpha_{\tau}^2/c^2\Phi_{f\tau}^{\mu} - \alpha\alpha_{\tau}/c^2\Phi_{f\xi\tau}^{\mu} - \left(\alpha^2/c^2 - (\beta_1^2 + \beta_2^2 + \beta_3^2)\right)\Phi_{f\xi\xi}^{\mu} = g_1\Phi_i\Phi_f^{\mu} + g_2\Phi_f^{\mu}.$$

After solving for any arbitrary parameters/coefficients and constants whenever possible and setting g_1 and g_2 to *unity*, then assuming the self and mixing interaction mass-energy equivalence equation, one can derive the following mass-energy equivalences:

$$m_{\psi_{\text{muon}}} = \frac{1}{72}\pi(-6 + \pi^2)|d1_{12}|^2,$$

$$m_{\psi_{\text{muon neutrino}}} = \frac{1}{288}\pi(-6 + \pi^2)|9\,\mathrm{M}^4 - 4\left(2b2_{12}{}^2 + d1_{12}{}^2\right)|,$$

$$m_{\phi_{pion}} = \frac{1}{32}\pi(-6 + \pi^2)|M|^4,$$

and

$$m_{\phi_f^{\mu}} = \frac{1}{72}\pi(-6 + \pi^2)|b2_{12}|^2.$$

After using the three most former expressions and the known rest masses for the particles or relativistic strings of interest, one can solve for the arbitrary parameters/coefficients $b2_{12}$, $d1_{12}$, and constant M using known values for items of interest. Next, the rest mass of the unknown particle or relativistic string is calculated to be 1.70 *X* 10^7 eV or 17 MeV.

4.2.) Glueball estimations

A glueball is a hypothetical particle solely comprised of gluons [58]. Particle physics theories suggest contemporary colliders should detect them. Although there is anecdotal evidence pointing to the existence of these particles, they have not been explicitly identified [47,58].

Since a gluon is a vector boson ϕ_i^{μ} and the Lagrangian terms for interactions of the entanglement type are as follows:

$$L_{interactions}[\phi_f, \phi_i^{\mu}, \psi_{i1}, \psi_{i2}] = g_1\left(\left|\phi_i^{\mu}\psi_{i1}\right|\phi_f + \left|\phi_i^{\mu}\psi_{i2}\right|\phi_f\right),$$

then, the principle of least action for the decay of a meson should be as follows:

$$S[\phi_f, \phi_i^{\mu}, \psi_{i1}, \psi_{i2}] = \int dx^4 \begin{pmatrix} -g_1\left(|\phi_i^{\mu}\psi_{i1}|\phi_f + |\phi_i^{\mu}\psi_{i2}|\phi_f\right) - g_2\left(|\psi_{i1}|^2 + |\psi_{i2}|^2 + 2|\phi_i^{\mu}|^2\right) \\ -\partial_\mu\phi_i^{\mu}\,\partial^\mu\phi_i'^{\mu} - \partial_\mu\psi_{i1}\,\partial^\mu\psi_{i1}' - \partial_\mu\psi_{i2}\,\partial^\mu\psi_{i2}' \\ -\frac{1}{2}c^2M^2\phi_f\phi_f' + \frac{1}{2}\partial_\mu\phi_f\,\partial^\mu\phi_f' \end{pmatrix}.$$

Suggesting ϕ_f is the mesonic effective or string field, Φ_i is the gluonic effective or string field, and Ψ_{i1} is one of the valence quark effective or string fields, while Ψ_{i2} is the other valence quark effective or string field present in the decay of the meson, the functional derivative of the principle of least action yields at least four transformed Hamiltonian equations:

$$\alpha i\Psi_{i1\xi} + \alpha_\tau i\Psi_{i1\tau} - \alpha_\tau^2/c^2\Psi_{i1\tau\tau} - \alpha\alpha_\tau/c^2\Psi_{i1\xi\tau} - \left(\alpha^2/c^2 - (\beta_1^2 + \beta_2^2 + \beta_3^2)\right)\Psi_{i1\xi\xi} = g_1\Phi_f\Psi_{i1} + g_2\Psi_{i1},$$

$$\alpha i\Psi_{i2\xi} + \alpha_\tau i\Psi_{i2\tau} - \alpha_\tau^2/c^2\Psi_{i2\tau\tau} - \alpha\alpha_\tau/c^2\Psi_{i2\xi\tau} - \left(\alpha^2/c^2 - (\beta_1^2 + \beta_2^2 + \beta_3^2)\right)\Psi_{i2\xi\xi} = g_1\Phi_f\Psi_{i2} + g_2\Psi_{i2},$$

$$\left(\alpha^2/c^2 - (\beta_1^2 + \beta_2^2 + \beta_3^2)\right)\Phi_{f\xi\xi} + c^2M^2\Phi_i = g_1\left(\Phi_i'^{\mu}\Psi_{i1} + \Psi_{i2}'\Phi_i^{\mu}\right),$$

and

$$\alpha i\Phi_{i\xi}^{\mu} + \alpha_\tau i\Phi_{i\tau}^{\mu} - \alpha_\tau^2/c^2\Phi_{i\tau}^{\mu} - \alpha\alpha_\tau/c^2\Phi_{i\xi\tau}^{\mu} - \left(\alpha^2/c^2 - (\beta_1^2 + \beta_2^2 + \beta_3^2)\right)\Phi_{i\xi\xi}^{\mu} = g_1\Phi_f\Phi_i^{\mu} + g_2\Phi_i^{\mu}.$$

After solving for any arbitrary parameters/coefficients and constants whenever possible and setting g_1 and g_2 to *unity*, and using the self and mixing interaction mass-energy equivalence equation, one can derive the following mass-energy equivalences:

$$m_{\psi_{i1}} = \frac{1}{72}\pi(-6 + \pi^2)|d1_{12}|^2,$$

$$m_{\psi_{i2}} = \frac{\pi(-6+\pi^2)\left|\frac{(9\,M^4 - 4b2_{12}d1_{12})^2}{b2_{12}^{\ 2}}\right|}{1152},$$

$$m_{\phi_f} = \frac{1}{32}\pi(-6 + \pi^2)|M|^4,$$

and

$$m_{\phi_i^{\mu}} = \frac{1}{72}\pi(-6 + \pi^2)|b2_{12}|^2.$$

One may also assume that the gluonic effective or string field ϕ_i^{μ} forms a glueball. In other words, the mass of a glueball would constitute the mass-energy equivalence of gluonic effective or string field. After using the three most former expressions and the known rest masses for the particles or relativistic strings of interest, one can solve for the arbitrary parameters/coefficients $b2_{12}$, $d1_{12}$, and constant M using known values for items of interest. Check the table for calculations of glueballs.

A table of predicted glueball masses derived from various meson decays is featured below:

	Quark 1/$m_{\psi_{i1}}$ (eV)	Quark 2/$m_{\psi_{i2}}$ (eV)	Meson/m_{ϕ_f} (eV)	Glueball/$m_{\phi_i^{\mu}}$ (eV)
charged pion	$2.20*10^7$	$4.70*10^6$	$1.40*10^8$	$1.49*10^9$ [65]
neutral pion (pair)	$2.20*10^6$	$4.70*10^6$	$1.35*10^8$	$1.39*10^9$ [65]
neutral kaon	$4.70*10^6$	$9.60*10^7$	$4.98*10^8$	$1.72*10^9$ [65]
J/psi meson	$1.28*10^9$	$1.28*10^9$	$3.10*10^9$	$1.87*10^9$ [66,67]

4.3.) Assessment of tetraquarks mass via decay by-products.

Recently, LHC at CERN found that the four charmed tetraquark $c\bar{c}c\bar{c}$, which had a rest mass of 6.9 GeV, decayed into the vector meson and sigma glueball σ [98,99,100].

Assuming the tetraquark is a scalar bosonic entity ϕ_i and the Lagrangian term for entanglement interaction is as follows:

$$L_{interactions}[\phi_i, \phi_{f1}^{\mu}, \phi_{f2}^{\mu}] = 2g_1|\phi_{f1}^{\mu}\phi_{f2}^{\mu}|\phi_i,$$

then the principle of least action for $c\bar{c}c\bar{c}$ decay into vector *J/y* mesons and s glueball is defined as:

$$S[\phi_i, \phi_{f1}^{\mu}, \phi_{f2}^{\mu}] = \int dx^4 \begin{pmatrix} -2g_1|\phi_{f1}^{\mu}\phi_{f2}^{\mu}|\phi_i - 2g_2\left(|\phi_{f1}^{\mu}|^2 + |\phi_{f2}^{\mu}|^2\right) - \partial_\mu\phi_{f1}^{\mu}\,\partial^\mu\phi_{f1}'^{\mu} \\ -\partial_\mu\phi_{f2}^{\mu}\,\partial^\mu\phi_{f2}^{\mu} - \frac{1}{2}c^2M^2\phi_i\phi_i' + \frac{1}{2}\partial_\mu\phi_i\,\partial^\mu\phi_i' \end{pmatrix}.$$

One may also deem that the tetraquark effective or string field ϕ_i is comprised of the vector meson J/ψ or glueball σ effective or string fields ϕ_{f1}^{μ} and ϕ_{f2}^{μ}, respectively. The functional derivatives of the Lagrangian yields three evolution equations:

$$\left(\alpha^2/c^2 - (\beta_1^2 + \beta_2^2 + \beta_3^2)\right)\Phi_{i\xi\xi} + c^2M^2\Phi_i = g_1\left(\Phi_{f2}^{\mu}\Phi_{f1}'^{\mu} + \Phi_{f1}^{\mu}\Phi_{f2}'^{\mu}\right),$$

$$\alpha i\Phi_{f1\xi}^{\mu} + \alpha_\tau i\Phi_{f1\tau}^{\mu} - \alpha_\tau^2/c^2\Phi_{f1\tau\tau}^{\mu} - \alpha\alpha_\tau/c^2\Phi_{f1\xi\tau}^{\mu} - \left(\alpha^2/c^2 - (\beta_1^2 + \beta_2^2 + \beta_3^2)\right)\Phi_{f1\xi\xi}^{\mu} = g_1\Phi_i\Phi_{f1}^{\mu} + g_2\Phi_{f1}^{\mu},$$

and

$$\alpha i\Phi_{f2\xi}^{\mu} + \alpha_\tau i\Phi_{f2\tau}^{\mu} - \alpha_\tau^2/c^2\Phi_{f2\tau\tau}^{\mu} - \alpha\alpha_\tau/c^2\Phi_{f2\xi\tau}^{\mu} - \left(\alpha^2/c^2 - (\beta_1^2 + \beta_2^2 + \beta_3^2)\right)\Phi_{f2\xi\xi}^{\mu} = g_1\Phi_i\Phi_{f2}^{\mu} + g_2\Phi_{f2}^{\mu}.$$

After solving for any arbitrary parameters/coefficients and constants whenever possible and setting g_1 and g_2 to *unity*, then assuming the self and mixing interaction mass-energy equivalence equation, one can derive the following mass-energy equivalences:

$$m_{c\bar{c}c\bar{c}} = \frac{1}{32}\pi(-6 + \pi^2)|M|^4,$$

$$m_{J/\psi} = \frac{1}{72}\pi(-6 + \pi^2)|b2_{12}|^2,$$

and

$$m_\sigma = \frac{9\pi(-6+\pi^2)|\frac{M^8}{b2_{12}{}^2}|}{2048}.$$

After using the two latter expressions and the known rest masses for the particles or relativistic strings of interest, one can solve for the arbitrary parameters/coefficients $b2_{12}$ and constant M using known values for the vector meson J/ψ or glueball σ. Plugging in the value M into the former expression claims, the mass of the four-charm tetraquark is $6.90*10^9$ eV.

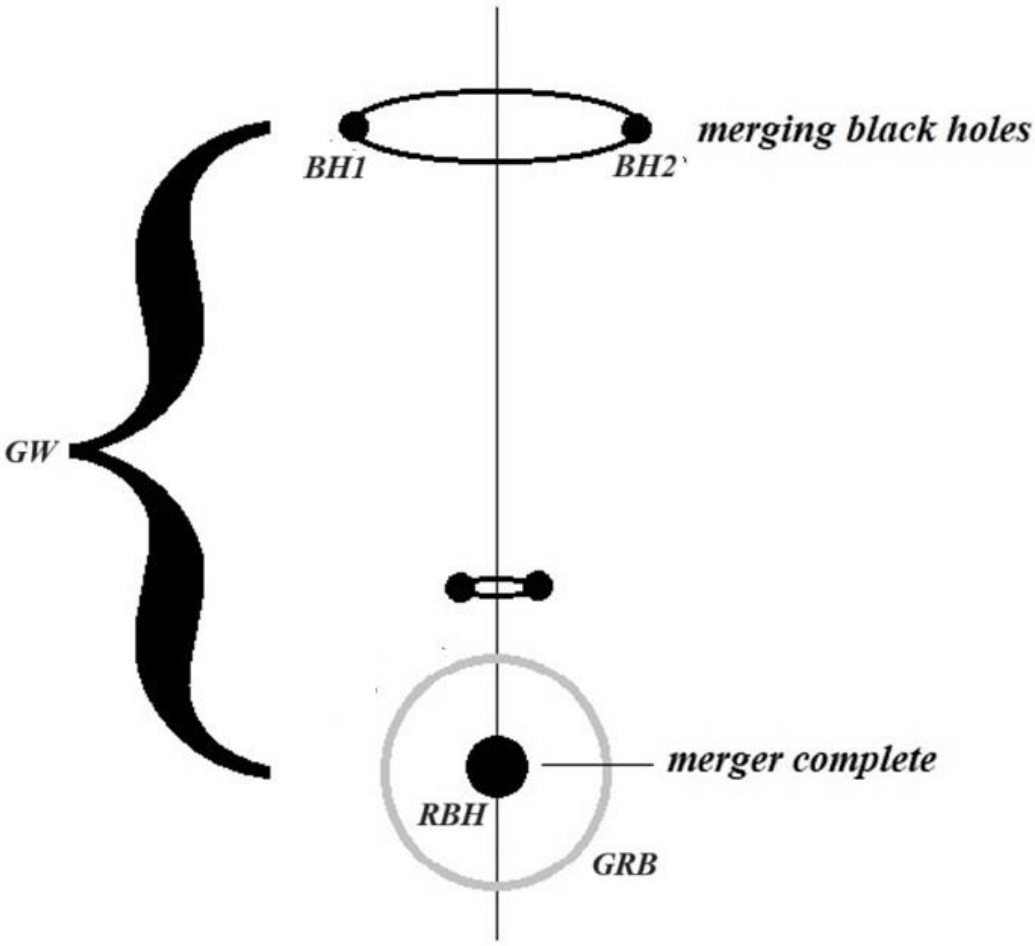
4.4.) Photon analysis linked to a binary black hole merger

Black holes (BH) were objects in the universe that had much speculation and excellent study. In the 18th century, Pierre-Simon Laplace and John Michell proposed that the universe possessed objects whose gravitational fields were so intense that light could not escape them [31]. Two centuries later, David Finkelstein and Karl Schwarzschild could make primitive solutions that were used to define such entities [69,70,93] initially. It was not until the 1960s that BH became a regular prediction in the general theory of relativity [37].

While undergoing mergers in spacetime, BH was known to emit gravitational waves (GW) and possibly some forms of light [49,94]. In 1916, Albert Einstein postulated that BH created ripples in spacetime in his great work, the General Theory of Relativity [46]. Before Einstein pinpointed the primary source of GW, Poincare, and Heaviside stated there were gravity's equivalent to electromagnetic waves [71,72]. In 2015, the first GW was detected by LIGO gravitational wave detectors. Finally, Atura Tanikawa and associates claimed that *g*-ray bursts were also emitted during the BH merger [74].

Quantum entanglement (QE) may be the primary mechanism for merging BHs. Leonard Susskind and Juan Maldacena generated a conjecture stating BH were like two entangled particles, or Einstein-Podolsky-Rosen pair, connected by a wormhole or Einstein-Rosen bridge [13]. Thus, they established the following relationship: ER = EPR. For this study, we used the previous equation to signify BH-QE.

In the first section of this study, we generate a Lagrangian to define BH-QE. Upon functional differentiation of this Lagrangian, we derive three quantum telegraph equations and two inhomogeneous nonlinear Klein-Gordon equations, or QT-KG. Assuming the three quantum telegraph equations described the two BH associated via ER = EPR and the merged BH while the inhomogeneous Klein-Gordon equations represented GW and photons, we used the generating function technique (GFT) to solve for the three BH, GW, and photons. Then, we try to predict the mass equivalents for photon emission, given we know the values of the three BH and GW. Ultimately, we concluded that population III stars were the source of BH mergers since they produced *g*-ray photons predicted via BH-QE system of equations.



Assuming black holes are vector boson-like and the Lagrangian term for interactions, such as entanglement, is as follows:

$$L_{interactions}[\phi_{i1}^{\mu}, \phi_{i2}^{\mu}, \phi_{f}^{\mu}, \phi_{G}^{\mu v}, \phi_{\gamma}^{\mu}] = g_1\phi_G^{\mu v}\left(|\phi_{i1}^{\mu}\phi_{\gamma}^{\mu}| + |\phi_{i2}^{\mu}\phi_{\gamma}^{\mu}| - |\phi_{f}^{\mu}|^2\right) g_{\mu v},$$

then the principle of least action for the merger of a binary black hole system is as follows:

$$S[\phi_{i1}^{\mu}, \phi_{i2}^{\mu}, \phi_{f}^{\mu}, \phi_{G}^{\mu v}, \phi_{\gamma}^{\mu}] =$$

$$\int dx^4 \left(\begin{array}{c} g_1\phi_G^{\mu v}\left(|\phi_{i1}^{\mu}\phi_{\gamma}^{\mu}| + |\phi_{i2}^{\mu}\phi_{\gamma}^{\mu}| - |\phi_{f}^{\mu}|^2\right) g_{\mu v} + g_2\left(|\phi_{i1}^{\mu}| + |\phi_{i2}^{\mu}| + 2|\phi_{\gamma}^{\mu}| - |\phi_{f}^{\mu}|\right) g_{\mu v} \\ + g_{\mu v}\,\partial_\mu\phi_{i1}^{\mu}\,\partial^\mu\phi_{i1}'^{\mu} + g_{\mu v}\,\partial_\mu\phi_{i2}^{\mu}\,\partial^\mu\phi_{i2}'^{\mu} + g_{\mu v}\,\partial_\mu\phi_{\gamma}^{\mu}\,\partial^\mu\phi_{\gamma}'^{\mu} - \phi_{f}^{\mu}\,\partial^\mu\phi_{f}'^{\mu} \\ -\frac{1}{2}c^2M^2\phi_G^{\mu v}\phi_G'^{\mu v}g_{\mu v} + \frac{1}{2}g_{\mu v}\,\partial_\mu\phi_G^{\mu v}\,\partial^\mu\phi_G'^{\mu v} + \partial_\mu\phi_G^{\mu v}\,\partial_v\phi_G'^{\mu v} + \partial_v\phi_G^{\mu v}\,\partial_\mu\phi_G'^{\mu v} \end{array} \right).$$

Deeming $\phi_G^{\mu v}$ is the gravitational wave/gravitonic practical or string field involved in the merger, ϕ_γ is the photonic effective or string field, ϕ_{i1}^{μ} is the first black hole, and ϕ_{i2}^{μ} is the second black hole effective or string fields involved in the merger, and ϕ_f^{μ} is the residual black hole effective or string field, the functional derivative of the principle of least action yields at least six transformed Hamiltonian equations:

$$\alpha i\Phi_{i1\xi}^{\mu} + \alpha_\tau i\Phi_{i1\tau}^{\mu} - \alpha_\tau^2/c^2\Phi_{i1\tau\tau}^{\mu} - \alpha\alpha_\tau/c^2\Phi_{i1\xi\tau}^{\mu} - \left(\alpha^2/c^2 - (\beta_1^2+\beta_2^2+\beta_3^2)\right)\Phi_{i1\xi\xi}^{\mu} = g_1\Phi_G^{\mu v}\Phi_{i1}^{\mu} + g_2\Phi_{i1}^{\mu},$$

and

$$\alpha i\Phi_{i2\xi}^{\mu} + \alpha_\tau i\Phi_{i2\tau}^{\mu} - \alpha_\tau^2/c^2\Phi_{i2\tau\tau}^{\mu} - \alpha\alpha_\tau/c^2\Phi_{i2\xi\tau}^{\mu} - \left(\alpha^2/c^2 - (\beta_1^2+\beta_2^2+\beta_3^2)\right)\Phi_{i2\xi\xi}^{\mu} = g_1\Phi_G^{\mu v}\Phi_{i2}^{\mu} + g_2\Phi_{i2}^{\mu},$$

$$\alpha i\Phi_{f\xi}^{\mu} + \alpha_\tau i\Phi_{f\tau}^{\mu} - \alpha_\tau^2/c^2\Phi_{f\tau\tau}^{\mu} - \alpha\alpha_\tau/c^2\Phi_{f\xi\tau}^{\mu} - \left(\alpha^2/c^2 - (\beta_1^2+\beta_2^2+\beta_3^2)\right)\Phi_{f\xi\xi}^{\mu} = g_1\Phi_G^{\mu v}\Phi_{f}^{\mu} + g_2\Phi_{f}^{\mu},$$

$$\alpha i\Phi_{\gamma\xi}^{\mu} + \alpha_\tau i\Phi_{\gamma\tau}^{\mu} - \alpha_\tau^2/c^2\Phi_{\gamma\tau\tau}^{\mu} - \alpha\alpha_\tau/c^2\Phi_{\gamma\xi\tau}^{\mu} - \left(\alpha^2/c^2 - (\beta_1^2+\beta_2^2+\beta_3^2)\right)\Phi_{\gamma\xi\xi}^{\mu} = g_1\Phi_G^{\mu v}\Phi_{\gamma}^{\mu} + g_2\Phi_{\gamma}^{\mu},$$

And

$$4AT\,\Phi_{G\xi\xi}^{\mu v} + \left(\alpha^2/c^2 - (\beta_1^2+\beta_2^2+\beta_3^2)\right)\Phi_{G\xi\xi}^{\mu v}g_{\mu v} + c^2M^2\Phi_G^{\mu v}g_{\mu v} = g_1((\Phi_{i1}^{\mu}\Phi_{\gamma}'^{\mu} + \Phi_{\gamma}^{\mu}\Phi_{i2}'^{\mu}) - \Phi_f^{\mu}\Phi_f'^{\mu})g_{\mu v},$$

$$\text{where } AT = \begin{pmatrix} -\frac{\alpha^2}{c^2} & 0 & 0 & 0 \\ 0 & -\beta_1^2 & 0 & 0 \\ 0 & 0 & -\beta_2^2 & 0 \\ 0 & 0 & 0 & -\beta_3^2 \end{pmatrix} \text{ and } g_{\mu v} = \begin{pmatrix} g_{11} & 0 & 0 & 0 \\ 0 & g_{22} & 0 & 0 \\ 0 & 0 & g_{33} & 0 \\ 0 & 0 & 0 & g_{44} \end{pmatrix},$$

AT is the Ansatz transform tensor, and $g_{\mu v}$ is the metric tensor. After solving for any arbitrary parameters/coefficients whenever possible and constants, setting g_1 and g_2 to *unity*, and applying the mass-energy equivalence equation to one element of the transformed solutions, assume the following relationships for determining the mass-energy equivalents of fermion and bosonic effective or string fields associated with the BH-QE system of equations were true:

$$m_{\phi_{i1}^{\mu}} = \frac{1}{72}\pi(-6+\pi^2)|b3_{12}|^2$$

$$m_{\phi_{i2}^{\mu}} = \frac{1}{4608}\pi(-6+\pi^2)\left|\frac{(9\,\mathrm{m}^4 + 8b2_{12}b3_{12} - 8b5_{12}{}^2)^2}{b2_{12}{}^2}\right|,$$

$$m_{\phi_{f}^{\mu}} = \frac{1}{72}\pi(-6+\pi^2)|b5_{12}|^2$$

$$m_{\phi_{G}} = \frac{1}{128}\pi(-6+\pi^2)|M|^4,$$

and

$$m_{\phi_{\gamma}^{\mu}} = \frac{1}{72}\pi(-6+\pi^2)|b2_{12}|^2$$

After using the four most former expressions and the known rest masses for the large-scale structures or relativistic strings of interest, one can solve for the arbitrary parameters/coefficients $b2_{12}$, $b3_{12}$, $b5_{12}$ and constant M using known values for items of interest. Ultimately, (s)he obtained the following table of photons for the first several GWs detected by LIGO:

Specific GW	BH 1/$m_{\phi_{i1}}$ (solar masses)	BH 2/$m_{\phi_{i2}}$ (solar masses)	Residual BH/m_{ϕ_f} (solar masses)	GW/m_{ϕ_G} (solar masses)	Photons/$m_{\phi_{\gamma}^{\mu}}$ (erg)
GW150914	35.6	30.6	63.1	3.1	$4.40*10^{55}$
GW151012	23.3	13.6	35.7	1.5	$2.60*10^{55}$
GW151226	13.7	7.7	20.5	1.0	$1.47*10^{55}$
GW170104	31.0	20.1	49.1	2.2	$3.54*10^{55}$
GW170608	10.9	7.6	17.8	0.9	$1.25*10^{55}$
GW170729	50.6	34.3	80.3	4.8	$5.26*10^{55}$
GW170809	35.2	23.8	56.4	2.7	$3.99*10^{55}$
GW170814	30.7	25.3	53.4	2.6	$3.69*10^{55}$
GW170818	35.4	26.7	59.4	2.7	$4.24*10^{55}$
GW170823	39.5	29.0	65.4	3.3	$4.56*10^{55}$

5.) Conclusion

Via *EFT* and *SFT*, one will likely identify particle X17 as the culprit causing the deviation in the muon g-2 experiment. The additional by-product from pion decay that coexists with the muons in the Fermilab storage ring has a rest mass of 17 MeV. This is the same mass as particle X17, a hypothetical protophobic spin-0 boson first captured by the ATOMKI and then JINR [101,102,103].

EFT and *SFT* might provide an accurate means to calculate glueballs derived from the decay of mesons. The rest of the masses of glueballs, condensed gluonic particles, or relativistic stings, established from the decay of pions and kaons, derived in this paper were consistent with CERN data [65]. Therefore, the theories discussed in this study could adequately estimate the masses of glueballs derived from the decay of other mesons.

EFT and *SFT* claimed that the initial gravitational waves detected by VIRGO/LIGO were generated from binary black hole mergers and were likely accompanied by g-ray bursts, or GRB. The theories in this study claimed the earliest observed binary black hole mergers likely emitted g-ray bursts at ~ 10^{55} erg. It is well known that many cosmological phenomena emit g-ray bursts. However, this range of energy for photons suggested the merging BH likely existed in the first generation or hypothetical population III stars [8].

Establishing a **Triality** between Quantum Information Theory (*QIT*), *EFT*, and *SFT*. Information theory is known as the study of uncertainty in the quantum realm, and its basic unit of information is the qubit, a two-state quantum mechanical system [104]. If an individual lets $c_2 = ic_1$, then (s)he obtains:

$$f(\xi) = c_1 e^{-\eta}$$

or

$$f(\xi) = c_1 e^{-i\xi}.$$

This expression is associated with the exponential map of a qubit [104]. Also, if one applies Euler's formula to the above expression, (s)he obtains the sinusoidal wave function f discussed in section 3. The above statements imply there is an intimate link between a qubit, physical body/particle, and string.

Conflicts of Interests

The author of this paper has no conflicts of interest.

References

1. Paspirou, "A BRIEF TOUR INTO THE HISTORY OF GRAVITY: FROM DEMOCRITUS TO EINSTEIN," American Journal of Space Science, vol. 1, no. 1, pp. 33–45, Jan. 2013, doi: 10.3844/ajssp.2013.33.45.
2. "A Geometric Theory of Everything," Scientific American, 2010.
3. R. L. Jackson, "A possible theory of partial differential equations," UVP, no. 3, 2021, doi: 10.21685/2072-3040-2021-3-3.
4. C. J.er and K. Rejzner, "Algebraic Quantum Field Theory -- an introduction." arXiv, Nov. 18, 2019. Accessed: Jan. 11, 2024. Online. Available: http://arxiv.org/abs/1904.04051
5. R. Penco, "An Introduction to Effective Field Theories".
6. K. Wray, "An Introduction to String Theory".
7. D. S. Stutts, "Analytical Dynamics: Lagrange's Equation and its Application – A Brief Introduction," THE CALCULUS OF VARIATIONS, 1995.
8. Y. Suwa and K. Ioka, "Gamma-ray bursts of the first generation of stars".
9. A. Blais, A. L. Grimsmo, S. M. Girvin, and A. Wallraff, "Circuit Quantum Electrodynamics," Rev. Mod. Phys., vol. 93, no. 2, p. 025005, May 2021, doi: 10.1103/RevModPhys.93.025005.
10. J. D. Martin, A. Roggero, H. Duan, J. Carlson, and V. Cirigliano, "Classical and quantum evolution in a simple coherent neutrino problem," Phys. Rev. D, vol. 105, no. 8, p. 083020, Apr. 2022, doi: 10.1103/PhysRevD.105.083020.
11. B. Zwiebach, "Closed String Field Theory: An Introduction." arXiv, May 07, 1993. Accessed: Jan. 11, 2024. Online. Available: http://arxiv.org/abs/hep-th/9305026
12. C. Y. Cardall, "Coherence of neutrino flavor mixing in quantum field theory," Phys. Rev. D, vol. 61, no. 7, p. 073006, Mar. 2000, doi: 10.1103/PhysRevD.61.073006.
13. J. Maldacena and L. Susskind, "Cool horizons for entangled black holes," Fortschritte der Physik, vol. 61, no. 9, pp. 781–811, Sep. 2013, doi: 10.1002/prop.201300020.
14. R. Durrer, "Cosmological Perturbation Theory".
15. P. Peter, "Cosmological Perturbation Theory." arXiv, Mar. 13, 2013. Accessed: Jan. 11, 2024. Online. Available: http://arxiv.org/abs/1303.2509
16. A. N. Ikot, O. A. Awoga, and B. I. Ita, "Exact Solutions of the Klein–Gordon Equation with Hylleraas Potential," Few-Body Syst, vol. 53, no. 3–4, pp. 539–548, Oct. 2012, doi: 10.1007/s00601-012-0434-y.
17. T. Gorringe, "Fermilab muon g -2 experiment," EPJ Web Conf., vol. 179, p. 01004, 2018, doi: 10.1051/epjconf/201817901004.
18. A. Kumar, "Fundamentals of Quantum Mechanics".
19. P. Fleury, Gravitation: from Newton to Einstein. 2019. doi: 10.1007/978-3-030-32001-0.
20. S. Scherer, "Introduction to Chiral Perturbation Theory." arXiv, Oct. 29, 2002. Accessed: Jan. 11, 2024. Online. Available: http://arxiv.org/abs/hep-ph/0210398
21. A. V. Manohar, "Introduction to Effective Field Theories." arXiv, Apr. 16, 2018. Accessed: Jan. 10, 2024. Online. Available: http://arxiv.org/abs/1804.05863
22. Y. Grossman, "Introduction to Flavour Physics," in LHC Phenomenology, E. Gardi, N. Glover, and A. Robson, Eds., Cham: Springer International Publishing, 2015, pp. 35–80. doi: 10.1007/978-3-319-05362-2_2.
23. P. Skands, "Introduction to QCD," in Searching for New Physics at Small and Large Scales, Nov. 2013, pp. 341–420. doi: 10.1142/9789814525220_0008.
24. J. Cardy, "Introduction to Quantum Field Theory".

25. D. Morin, “Introduction to quantum mechanics,” . INTRODUCTION TO QUANTUM MECHANICS.

26. W. Siegel, Introduction to string field theory. in Advanced series in mathematical physics, no. 8. Singapore: World Scientific, 1988.

27. R. Jackiw, “Introduction to the Yang-Mills quantum theory,” Rev. Mod. Phys., vol. 52, no. 4, pp. 661–673, Oct. 1980, doi: 10.1103/RevModPhys.52.661.

28. L. Álvarez-Gaumé and M. A. Vázquez-Mozo, “Introductory Lectures on Quantum Field Theory”.

29. R. H. Brandenberger, Lectures on the Theory of Cosmological Perturbations, vol. 646. 2004. doi: 10.1007/b97189.

30. M. Rinaldi and V. Vento, “Meson and glueball spectroscopy within the graviton soft wall model,” Phys. Rev. D, vol. 104, no. 3, p. 034016, Aug. 2021, doi: 10.1103/PhysRevD.104.034016.

31. C. Montgomery, W. Orchiston, and I. Whittingham, “MICHELL, LAPLACE AND THE ORIGIN OF THE BLACK HOLE CONCEPT,” J. Astron. Hist. Herit., vol. 12, no. 2, pp. 90–96, Jul. 2009, doi: 10.3724/SP.J.1440-2807.2009.02.01.

32. D. Cremer, “Møller–Plesset perturbation theory: from small molecule methods to methods for thousands of atoms,” WIREs Comput Mol Sci, vol. 1, no. 4, pp. 509–530, Jul. 2011, doi: 10.1002/wcms.58.

33. A. Keshavarzi, K. S. Khaw, and T. Yoshioka, “Muon \$g-2\$: A review,” Nuclear Physics B, vol. 975, p. 115675, Feb. 2022, doi: 10.1016/j.nuclphysb.2022.115675.

34. B. D. S. L. Torres, T. R. Perche, A. G. S. Landulfo, and G. E. A. Matsas, “Neutrino flavor oscillations without flavor states,” Phys. Rev. D, vol. 102, no. 9, p. 093003, Nov. 2020, doi: 10.1103/PhysRevD.102.093003.

35. BESIII Collaboration et al., “Observation of a structure at 1.84 GeV/c\$^2\$ in the \$3(\pi^+\pi^-)\$ mass spectrum in \$J/\psi\rightarrow \gamma 3(\pi^+\pi^-)\$ decays,” Phys. Rev. D, vol. 88, no. 9, p. 091502, Nov. 2013, doi: 10.1103/PhysRevD.88.091502.

36. B. P. Abbott et al., “Observation of Gravitational Waves from a Binary Black Hole Merger,” Phys. Rev. Lett., vol. 116, no. 6, p. 061102, Feb. 2016, doi: 10.1103/PhysRevLett.116.061102.

37. K. Schwarzschild, “On the gravitational field of a mass point according to Einstein’s theory.” arXiv, May 12, 1999. Accessed: Jan. 11, 2024. Online. Available: http://arxiv.org/abs/physics/9905030

38. E. P. Verlinde, “On the Origin of Gravity and the Laws of Newton,” J. High Energ. Phys., vol. 2011, no. 4, p. 29, Apr. 2011, doi: 10.1007/JHEP04(2011)029.

39. “(PDF) Cygnus X-1: A Spinning Black Hole?” Accessed: Jan. 11, 2024. Online. Available: https://www.researchgate.net/publication/1817344_Cygnus_X-1_A_Spinning_Black_Hole

40. D. M. Gingrich, “Practical Quantum Electrodynamics”.

41. A. S. Blum and D. Rickles, Eds., Quantum gravity in the first half of the twentieth century: a sourcebook = Edition Open Sources, 1. Auflage. in Edition Open Sources, no. 10. Berlin: Edition Open Access, 2018.

42. D. Rickles, “Quantum Gravity: A Primer for Philosophers”.

43. A. Westphal and D. Hamburg, “Quantum Mechanics and Gravitation”.

44. A. I. Arbab, “Quantum Telegraph equation: New matter wave equation,” Optik, vol. 140, pp. 1010–1019, Jul. 2017, doi: 10.1016/j.ijleo.2017.05.002.

45. B. S. DeWitt, “Quantum Theory of Gravity. I. The Canonical Theory,” Phys. Rev., vol. 160, no. 5, pp. 1113–1148, Aug. 1967, doi: 10.1103/PhysRev.160.1113.

46. A. Einstein, “Relativity: The Special and General Theory”.

47. A. S. Kronfeld and C. Quigg, “Resource Letter QCD-1: Quantum chromodynamics,” American Journal of Physics, vol. 78, no. 11, pp. 1081–1116, Nov. 2010, doi: 10.1119/1.3454865.

48. B. Moziak, “SEARCH FOR INVISIBLE DECAYS OF THE J/ψ RESONANCE”.

49. R. Perna, D. Lazzati, and B. Giacomazzo, “Short Gamma-Ray Bursts from the Merger of Two Black Holes,” ApJL, vol. 821, no. 1, p. L18, Apr. 2016, doi: 10.3847/2041-8205/821/1/L18.

50. I. Miškovičová, M. Hanke, J. Wilms, M. Nowak, K. Pottschmidt, and N. Schulz, “Spectroscopy of the Stellar Wind in the Cygnus X-1 System,” Acta Polytechnica,+ vol. 51, Mar. 2011, doi: 10.14311/1332.

51. P. Carenza, “Stockholm University, OKC”.

52. H. j, String Field Theory -- A Modern Introduction, vol. 980. 2021. doi: 10.1007/978-3-030-65321-7.

53. K. Becker, M. Becker, and J. H. Schwarz, “String Theory and M-Theory”.

54. M. H. Poincaré, “Sur la dynamique de l’électron,” Rend. Circ. Matem. Palermo, vol. 21, no. 1, pp. 129–175, Dec. 1906, doi: 10.1007/BF03013466.

55. J. Mott, “The Muon g − 2 experiment at Fermilab”.

56. S. Godfrey, “THE PHENOMENOLOGY OF GLUEBALL AND HYBRID MESONS”.

57. H. Zinkernagel, “The Philosophy behind Quantum Gravity,” THEORIA, vol. 21, no. 3, pp. 295–312, Sep. 2006, doi: 10.1387/theoria.522.

58. V. Mathieu, N. Kochelev, and V. Vento, “The Physics of Glueballs,” Int. J. Mod. Phys. E, vol. 18, no. 01, pp. 1–49, Jan. 2009, doi: 10.1142/S0218301309012124.

59. J. Binney, “The Physics of Quantum Mechanics”.

60. J. D. Fraser, “The Real Problem with Perturbative Quantum Field Theory,” The British Journal for the Philosophy of Science, vol. 71, no. 2, pp. 391–413, Jun. 2020, doi: 10.1093/bjps/axx042.

61. R. B. Laughlin and D. Pines, “The Theory of Everything,” Proc. Natl. Acad. Sci. U.S.A., vol. 97, no. 1, pp. 28–31, Jan. 2000, doi: 10.1073/pnas.97.1.28.

62. C. Leung, “THE THEORY OF EVERYTHING”.

63. R. L. Jackson, “The Utilization of the Generating Function Technique in the Discovery of Solutions

for the Three-Dimensional Navier-Stokes Equation System," OJFD, vol. 12, no. 01, pp. 86–95, 2022, doi: 10.4236/ojfd.2022.121005.

64. C.K. Skylaris, "Pertubation Theory," Physical & Theoretical Chemistry Laboratory South Parks Road, Oxford 2006

65. C. Amsler, C. Hanhart, "Non-qq mesons," Particle Data Group Review. 2017

66. M. Ablikim, M.N. Achasov, O. Albayrak, D.J. Ambrose, F.F. An, Q. An, BESIII Collaboration, "Observation of a structure at 1.84 GeV/c(2) in the 3(pi(+)pi(-)) mass spectrum in J/psi -> gamma 3(pi(+) pi(-)) decays." Physical Review D, 88 (9). https://dx.doi.org/10.1103/PhysRevD.88.091502. 2013

67. B. Moziak, "Search for invisible decays of the J/psi resonance," Rensselaer Polytechnic Institute ProQuest Dissertations Publishing, 2010

68. Martinez-Nuñez, S., Kretschmar, P., Bozzo, E., Oskinova, L. M., Puls, J., Sidoli, L., Sundqvist, J. O., Blay, P., Falanga, M., Fürst, F., Giménez-García, A., Kreykenbohm, I., Kühnel, M., Sander, A., Torrejón, J. M., & Wilms, J. "Towards a unified view of inhomogeneous stellar winds in isolated supergiant stars and supergiant high mass X-ray binaries." 2018 https://doi.org/10.1007/s11214-017-0340-1

69. B. L. Webster, P. Murdin, "Cygnus X-1—a Spectroscopic Binary with a Heavy Companion?", Nature, 235 (5332): 37–38. 1972

70. K. Finkelstein, "Past-Future Asymmetry of the Gravitational Field of a Point Particle," Phys. Rev. 112 (May 15): 965-967 1958

71. H. Poincare, "Sur la dynamique de l' electron." C.R. T.140 (1905) 1504-1508

72. O. Heaviside, "A Gravitational and Electromagnetic Analogy," Part I, The Electrician, 31: 281-282 1983

73. B.P. Abbott, (LIGO Scientific Collaboration and Virgo Collaboration). "Observation of Gravitational Waves from a Binary Black Hole Merger," Phys. Rev. Lett. 116 (6): 061102. 2016

74. R. Perna, D. Lazzati, B. Giacomazzo, "Short Gamma-Ray Bursts from he Merger of Two Black Holes," The Astrophysical Lett. 821:L18 (6pp), 2016 April 10: 1-6. 2016

75. H. Colakoglu, "Volume of a Tetrahedron in the Taxicab Space". In: Missouri Journal of Mathematical Sciences 21, pp. 21–27. 2009

76. K. Wray, "An Introduction to String Theory".

77. R. C. Myers, "Nonabelian Phenomena on D-branes," Class. Quantum Grav., vol. 20, no. 12, pp. S347–S372, Jun. 2003, doi: 10.1088/0264-9381/20/12/302.

78. G. W. Moore and G. Segal, "D-branes and K-theory in 2D topological field theory".

79. G. Moore, N. Nekrasov, and S. Shatashvili, "D -Particle Bound States and Generalized Instantons," Communications in Mathematical Physics, vol. 209, no. 1, pp. 77–95, Jan. 2000, doi: 10.1007/s002200050016.

81. D. S. Freed and E. Witten, "Anomalies in string theory with D-branes," Asian Journal of Mathematics, vol. 3, no. 4, pp. 819–852, 1999, doi: 10.4310/AJM.1999.v3.n4.a6.

82. M. R. Douglas, D. Kabat, P. Pouliot, and S. H. Shenker, "D-branes and Short Distances in String Theory," Nuclear Physics B, vol. 485, no. 1–2, pp. 85–127, Feb. 1997, doi: 10.1016/S0550-3213(96)00619-0.

83. R. Dijkgraaf, L. Hollands, P. Sułkowski, and C. Vafa, "Supersymmetric gauge theories, intersecting branes and free fermions," J. High Energy Phys., vol. 2008, no. 02, pp. 106–106, Feb. 2008, doi: 10.1088/1126-6708/2008/02/106.

84. R. Blumenhagen, M. Cvetiˇc, S. Kachru, and T. Weigand, "D-brane Instantons in Type II String Theory".

85. E. Bergshoeff, J. Lahnsteiner, L. Romano, and J. Rosseel, "The Supersymmetric Neveu-Schwarz Branes of Non-Relativistic String Theory," J. High Energ. Phys., vol. 2022, no. 8, p. 218, Aug. 2022, doi: 10.1007/JHEP08(2022)218.

86. D. Berenstein and E. Dzienkowski, "Open spin chains for giant gravitons and relativity," J. High Energ. Phys., vol. 2013, no. 8, p. 47, Aug. 2013, doi: 10.1007/JHEP08(2013)047.

87. V. Balasubramanian, A. Kar, S. F. Ross, and T. Ugajin, "Spin structures and baby universes," J. High Energ. Phys., vol. 2020, no. 9, p. 192, Sep. 2020, doi: 10.1007/JHEP09(2020)192.

88. Hajarolasvadi, S. "A new hybrid numerical scheme for simulating fault ruptures with near fault bulk inhomogeneities." 2016 https://core.ac.uk/download/158315427.pdf

89. 이., & 조. 근사 원방 경계조건을 이용한 HAWT의 Navier-Stokes 유동해석. 한국항공우주학회지. 2008 http://www.dbpia.co.kr/Article/932962

90. Electroweak epoch - Wikipedia. https://en.wikipedia.org/wiki/Electroweak_epoch

91. What is astrophysics? | Space | EarthSky. https://earthsky.org/space/definition-what-is-astrophysics/

92. Redhead, M. (1980). Models in Physics. The British Journal for the Philosophy of Science. https://doi.org/10.1093/bjps/31.2.145

93. S. Sadiq, A. Alias. "Black holes formations and dark fabric distortions." 2022 https://core.ac.uk/download/524625565.pdf

94. B. Zhang, B. "The Delay Time of Gravitational Wave — Gamma-Ray Burst Associations." 2019 https://core.ac.uk/download/225356190.pdf

95. M. M. Ivanov, "Effective Field Theory for Large Scale Structure." arXiv, Dec. 16, 2022. Accessed: Jan. 26, 2024. Online. Available: http://arxiv.org/abs/2212.08488

96. Z. Vlah, U. Seljak, M. Y. Chu, and Y. Feng, "Perturbation theory, effective field theory, and oscillations in the power spectrum," J. Cosmol. Astropart. Phys., vol. 2016, no. 03, pp. 057–057, Mar. 2016, doi: 10.1088/1475-7516/2016/03/057.

97. H. Erbin, String Field Theory -- A Modern Introduction, vol. 980. 2021. doi: 10.1007/978-3-030-65321-7.

98. Z. Kuang, K. Serafin, X. Zhao, and J. P. Vary, "All-charm tetraquark in front form dynamics," Phys. Rev. D, vol. 105, no. 9, p. 094028, May 2022, doi:

10.1103/PhysRevD.105.094028.

99. R. J. Lloyd and J. P. Vary, “All-charm tetraquarks,” Phys. Rev. D, vol. 70, no. 1, p. 014009, Jul. 2004, doi: 10.1103/PhysRevD.70.014009.

100. R. Tiwari, D. P. Rathaud, and A. K. Rai, “Spectroscopy of all charm tetraquark states,” Indian J Phys, vol. 97, no. 3, pp. 943–954, Mar. 2023, doi: 10.1007/s12648-022-02427-8.

101. K. U. Abraamyan et al., “Observation of structures at ~ 17 and ~ 38 MeV/c^{2} in the $\gamma\gamma$ invariant mass spectra in pC, dC, and dCu collisions at $\textit{p}_{lab}$ of a few GeV/c per nucleon.” arXiv, Nov. 30, 2023. Accessed: Jan. 26, 2024. Online. Available: http://arxiv.org/abs/2311.18632

102. Y. Hiçyılmaz, S. Khalil, and S. Moretti, “Light $Z’$ Signatures at the LHC,” Phys. Rev. D, vol. 107, no. 3, p. 035030, Feb. 2023, doi: 10.1103/PhysRevD.107.035030.

103. D. S. Firak et al., “Confirmation of the existence of the X17 particle,” EPJ Web Conf., vol. 232, p. 04005, 2020, doi: 10.1051/epjconf/202023204005.

104. A. J. et al., “Quantum Algorithm Implementations for Beginners,” ACM Transactions on Quantum Computing, vol. 3, no. 4, pp. 1–92, Dec. 2022, doi: 10.1145/3517340.

105. J. Dunkel and P. Hänggi, “Relativistic Brownian motion,” Physics Reports, vol. 471, no. 1, pp. 1–73, Feb. 2009, doi: 10.1016/j.physrep.2008.12.001.